# Aliasing Free for Mixed Spectra for Stable Processes 

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#### Abstract

This work focuses on the symmetric alpha stable processes with continuous time frequently used in modeling the signal with indefinitely growing variance when the spectral measure is mixed: sum of a continuous meseare and discrete measure. The objective of this paper is to estimate the spectral density of the continuous part from discrete observations of the signal. For that, we propose a method based on a sample of the signal at a periodic instant. The Jackson polynomial kernel is used for construct a periodogram. We smooth this periodogram by two spectral windows taking into account the width of the interval where the spectral density is nonzero. This technique allows to circumvent the phenomenon of aliasing often encountered in the estimation from the discrete observations of a process with a continuous time.


## KEYWORDS

Spectral density, stable processes, periodogram, smoothing estimate, aliasing

## 1. INTRODUCTION

Stable alpha processes have been of interest to several research authors for their multiple applications when we have random signals with variance indefinitely increase. The harmonizable process is an important example of a symmetric $\alpha$-stable process, and its proprieties have been considered by numerous authors like [1]-[10] to name a few.

In particular, stable symmetric processes find their place in various applications and in various fields such as: physics, biology, electronics and electricity, hydrology, economies, communications and radar applications, ...ect. See: [11]-[22]. This work considers a symmetric alpha stable harmonizable process $X=\{X(t): t \in R\}$. Alternatively $X$ has the integral representation:

$$
\begin{equation*}
X(t)=\int \exp [i(t \lambda)] d \xi(\lambda) \tag{1}
\end{equation*}
$$

where $1<\alpha<2$ and ${ }^{\xi}$ is a complex valued symmetric $\alpha_{\text {-stable random measure on } R}$ with independent and isotropic increments. The measure defined by $m(A)=|\xi(A)|_{\alpha}^{\alpha}$ (see [4]) is called "control" measure or spectral measure." The spectral density function was already estimated in different cases: by E.Masry and S.Combanis [4] when the time of the process is continuous, by Sabre [23] when the time of the process is discrete and by R. Sabre [24]-[25] when the time of the process is p -adic.

This work considers a general case where we suppose that the spectral measure is the sum of an absolutely continuous measure with respect to Lebesgue measure and a discrete measure:

$$
d \mu(\lambda)=\phi(x) d x+\sum_{i=1}^{q} c_{i} \delta_{w_{i}}
$$

where $\delta$ is a Dirac measure, $\phi$ is nonnegative integrable and bounded function. $c_{i}$ is an unknown positive real number and $w_{i}$ is an unknown real number. Assume that $w_{i} \neq 0$. The function $\phi^{\phi}$ is called the spectral density. Discrete measure is due to random repeated value jumps during experimental measurements. The spectral density $\phi$ represents the distribution of the energy carried by the signal.

Our goal is to establish a non-parametric estimate of the spectral density $\phi$ from discrete observations of $X(t)$. This is motived by the fact that, in practice, it is not obvious to observe the process on continuous interval of time. Indeed, we sampled the process at instants $t_{n}$, equally distant, i.e., $t_{n}=n \tau, \tau>0$. It is known that aliasing of $\phi$ occurs. For more details about aliasing phenomenon, see [26]. To avoid this difficulty, we suppose that the spectral density $\phi$ is vanishing for $|\lambda|>\Omega$ where $\Omega$ is a nonnegative real number. From some smoothing, we construct an estimate depending on $\Omega$ and we show that it is asymptotically unbiased and consistent.

Briefly, we indicate the organisation of this paper: the outline in this paper is as follows: we present in the second section two technical lemmas, the preiodogram and show that this periodogram is asymptotically unbiased estimated but not consistent. In the third section, we smooth this periodogram by two chosen spectral windows to estimate the spectral density at jump points. We show that the smoothing periodogram is a consistent estimator.

## 2. The Periodogram and its Proprieties

First, we introduce some basic notations and properties of the Jackson's polynomial kernel. Let $N$ is the size of sample of $X$. Let $k$ and $n$ are the numbers satisfying: $N-1=2 k(n-1) \quad$ with $\quad n \in N \quad k \in N \cup\left\{\frac{1}{2}\right\} \quad k=\frac{1}{2} \quad$ if $\quad n=2 n_{1}-1, n_{1} \in N$.

The Jackson's polynomial kernel is defined by: $\left|H_{N}(\lambda)\right|^{\alpha}=\left|A_{N} H^{(N)}(\lambda)\right|^{\alpha} \quad$ where

$$
H^{(N)}(\lambda)=\frac{1}{q_{k, n}}\left(\frac{\sin \frac{n \lambda}{2}}{\sin \frac{\lambda}{2}}\right)^{2 k} \text { with } q_{k, n}=\frac{1}{2 \pi} \int_{-\pi}^{\pi}\left(\frac{\sin \frac{n \lambda}{2}}{\sin \frac{\lambda}{2}}\right)^{2 k} d \lambda .
$$

where $A_{N}=\left(B_{\alpha, N}\right)^{\frac{-1}{\alpha}}$ with $B_{\alpha, N}=\int_{-\pi}^{\pi}\left|H^{(N)}(\lambda)\right|^{\alpha} d \lambda$.
We give the following lemmas which are used in the rest of this paper. Their proof are given in [23].

Lemma 2.1 There is a non negative function $h_{k}$ such as:

$$
\begin{gathered}
H^{(N)}(\lambda)=\sum_{m=-k(n-1)}^{k(n-1)} h_{k}\left(\frac{m}{n}\right) \cos (m \lambda) \\
B_{\alpha, N}^{\prime}=\int_{-\pi}^{\pi}\left|\frac{\sin \frac{n \lambda}{2}}{\sin \frac{\lambda}{2}}\right|^{2 k \alpha} d \lambda \quad \text { and } \quad J_{N, \alpha}=\int_{-\pi}^{\pi}|u|^{\gamma}\left|H_{N}(u)\right|^{\alpha} d u, \\
\text { where }
\end{gathered}
$$ $\gamma \in] 0,2]$, then

$$
B_{\alpha, N}^{\prime}\left\{\begin{array}{lll}
\geq 2 \pi\left(\frac{2}{\pi}\right)^{2 k \alpha} n^{2 k \alpha-1} & \text { if } & 0<\alpha<2 \\
\leq \frac{4 \pi k \alpha}{2 k \alpha-1} n^{2 k \alpha-1} & \text { if } & \frac{1}{2 k}<\alpha<2
\end{array}\right.
$$

And

$$
J_{N, \alpha} \leq\left\{\begin{array}{lll}
\frac{\pi^{\gamma+2 k \alpha}}{2^{2 k \alpha}(\gamma-2 k \alpha+1)} \frac{1}{n^{2 k \alpha-1}} & \text { if } & \frac{1}{2 k}<\alpha<\frac{\gamma+1}{2 k} \\
\frac{2 k \alpha \pi^{\gamma+2 k \alpha}}{2^{2 k \alpha}(\gamma+1)(2 k \alpha-\gamma-1)} \frac{1}{n^{\gamma}} & \text { if } & \frac{\gamma+1}{2 k}<\alpha<2
\end{array}\right.
$$

In this section, we give a periodogram and we develop its proprieties. Assume that the process $X(t)$, defined in (1), is observed at instants $t_{j}=j \tau, j=1,2, \ldots N$ and $\tau=\frac{2 \pi}{\omega}$, where $\omega$ is a real number strictly greater than $2 \Omega$. We define the periodogram $\hat{I}_{N}$ on $]-\Omega, \Omega[$ as follows:

$$
\hat{I}_{N}(\lambda)=C_{p, \alpha}\left|I_{N}(\lambda)\right|^{p}, \quad 0<p<\frac{\alpha}{2}
$$

where

$$
I_{N}(\lambda)=[\tau]^{\frac{1}{\alpha}} A_{N} R \mathrm{e}\left[\sum_{n^{\prime}=-k(n-1)}^{n^{\prime}=k(n-1)} h_{k}\left(\frac{n^{\prime}}{n}\right) \exp \left\{-i\left(n^{\prime} \tau \lambda\right)\right\} X\left(n^{\prime} \tau+k(n-1) \tau\right)\right],
$$

and the normalisation constant $C_{p, \alpha}$ is given by $C_{p, \alpha}=\frac{D_{p}}{F_{p, \alpha}\left[C_{\alpha}\right]^{p / \alpha}}$, with

$$
D_{p}=\int_{-\infty}^{\infty} \frac{1-\cos (u)}{|u|^{1+p}} d u \quad \text { and } \quad F_{p, \alpha}=\int_{-\infty}^{\infty} \frac{1-e^{-|u|^{a}}}{|u|^{1+p}} d u
$$

## Lemma 2.3

The characteristic function of $\quad I_{N}(\lambda), \quad \operatorname{Eexp}\left[i r I_{N}(\lambda)\right]$, converges to $\exp \left[-C_{\alpha}|r|^{\alpha}\left(\psi_{N, 1}(\lambda)+\psi_{N, 2}(\lambda)\right]\right.$. where

$$
\psi_{N, 1}(\lambda)=\int_{-\pi}^{\pi}\left|H_{N}(y-\tau \lambda)\right|^{\alpha} \phi\left(\frac{y}{\tau}\right) d y \quad \psi_{N, 2}(\lambda)=\sum_{i=1}^{q} c_{i}\left|H_{N}\left(w_{i}-\tau \lambda\right)\right|^{\alpha}
$$

## Proof

By substituting (1) in the expression of $I_{N}$, we have:

$$
I_{N}(\lambda)=[\tau]^{\frac{1}{\alpha}} A_{N} R e \int_{R_{n^{\prime}=-k(n-1)}}^{n^{\prime}=k(n-1)} h_{k}\left(\frac{n^{\prime}}{n}\right) \exp \left\{i\left[n^{\prime} \tau(\lambda-u)\right]\right\} \exp \{i[\tau u k(n-1)]\} d \xi(u)
$$

It follows from [1] and the definition of the Jackson polynomial kernel that the characteristic function is the form:

$$
\begin{equation*}
E \exp \left[i r I_{N}(\lambda)\right]=\exp \left[-C_{a}|r|^{a} \psi_{N}(\lambda)\right] \tag{2}
\end{equation*}
$$

where $\psi_{N}(\lambda)=\psi_{N, 1}(\lambda)+\psi_{N, 2}(\lambda) \quad$ with
$\psi_{N, 1}(\lambda)=\int_{R}\left|H_{N}(v-\tau \lambda)\right|^{a} \phi\left(\frac{v}{\tau}\right) d v \quad$ and $\quad \psi_{N, 2}(\lambda)=\sum_{i=1}^{q} c_{i}\left|H_{N}\left(w_{i}-\tau \lambda\right)\right|^{\alpha}$
We can write $\psi_{N, 1}(\lambda)=\int_{R}\left|H_{N}(v-\tau \lambda)\right|^{a} \phi\left(\frac{v}{\tau}\right) d v=\sum_{j \in \mathcal{Z}} \int_{(2 j-1) \pi}^{(2 j+1) \pi}\left|H_{N}(v-\tau \lambda)\right|^{a} \phi\left(\frac{v}{\tau}\right) d v$.
Putting $v=y-2 \pi j$ and using the fact that $H_{N}$ is $2 \pi$-periodic, we obtain $\psi_{N}(\lambda)=\sum_{j \in Z} \int_{-\pi}^{\pi}\left|H_{N}(y-\tau \lambda)\right|^{a} \phi_{j}(y) d y,{ }_{\text {where }} \phi_{j}(y)=\phi\left(\frac{y}{\tau}-\frac{2 \pi}{\tau} j\right) .{ }_{\text {Let }} j$ be an integer
such that $-\Omega<\frac{y-2 \pi j}{\tau}<\Omega$. Using the fact that $\tau \Omega<\pi$ and $|y|<\pi$, we get $|j|<\frac{\tau \Omega}{2 \pi}+\frac{1}{2}<1$ and then $j=0$. Consequently:

$$
\begin{equation*}
\psi_{N, 1}(\lambda)=\int_{-\pi}^{\pi}\left|H_{N}(y-\tau \lambda)\right|^{\alpha} \phi\left(\frac{y}{\tau}\right) d y \tag{3}
\end{equation*}
$$

Theorem 2.4 Let $-\Omega<\lambda<\Omega$ then

$$
E\left[\hat{I}_{N}(\lambda)\right]=\left[\psi_{N}(\lambda)\right]^{\frac{p}{\alpha}}
$$

## Proof

As in [4], we use the following equality: for all real $x$ and $0<p<2$,

$$
\begin{equation*}
|x|^{p}=D_{p}^{-1} \int_{-\infty}^{\infty} \frac{1-\cos (x u)}{|u|^{1+p}} d u=D_{p}^{-1} R \mathrm{e} \int_{-\infty}^{\infty} \frac{1-e^{i x u}}{|u|^{1+p}} d u \tag{4}
\end{equation*}
$$

replacing ${ }^{x}$ by $I_{N}$, we obtain

$$
\begin{equation*}
\hat{I}_{N}(\lambda)=\frac{1}{F_{p, \alpha}\left[C_{a}\right]^{p / \alpha}} R \mathrm{e} \int_{-\infty}^{\infty} \frac{1-\exp \left\{i u I_{N}(\lambda)\right\}}{|u|^{1+p}} d u \tag{5}
\end{equation*}
$$

Using (6) and the definition of the $F_{p, \alpha}$, we get

$$
\begin{aligned}
E \hat{I}_{N}(\lambda) & =\frac{1}{F_{p, \alpha}\left[C_{a}\right]^{p / \alpha}} \int_{R} \frac{1-\exp \left\{-C_{a}|u|^{\alpha} \psi_{N}(\lambda)\right\}}{|u|^{1+p}} d u . \\
& =\left[\psi_{N}(\lambda)\right]^{p / \alpha}
\end{aligned}
$$

## 3. Smoothing Periodogram

In order to obtain a consistent estimate of $[\phi(\lambda)]^{\frac{p}{\alpha}}$, we smooth the periodogram via spectral windows depending on whether $\tau \lambda$ is a jump point or not $\left(\tau \lambda \neq w_{i}\right)$.

$$
f_{N}(\lambda)=\left\{\begin{array}{c}
f_{N}^{(1)}(\lambda) \text { if } \tau \lambda \notin\left\{w_{1}, w_{2}, \cdots, w_{q}\right\} \\
\frac{f_{N}^{(2)}(\lambda)-c f_{N}^{(1)}(\lambda)}{1-c} \text { else }
\end{array}\right.
$$

where $f_{N}^{(1)}(\lambda)=\int_{-\pi}^{\pi} W_{N}^{(1)}(\lambda-u) \hat{I}_{N}(u) d u$ and $f_{N}^{(2)}(\lambda)=\int_{-\pi}^{\pi} W_{N}^{(2)}(\lambda-u) \hat{I}_{N}(u) d u$.

The spectral windows $W_{N}^{(1)}$ and $W_{N}^{(2)}$ are defined by: $W_{N}^{(1)}(x)=M_{N}^{(1)} W\left(M_{N}^{(1)} x\right)$ and $W_{N}^{(2)}(x)=M_{N}^{(2))} W\left(M_{N}^{(2)} x\right)$ with $W$ is an even nonnegative, continuous function, vanishing for $|\lambda|>1$ such that $\int_{-1}^{1} W(u) d u=1$. The bandwidths $\left.M_{N}^{(1)}\right)$ and $\left.M_{N}^{(2)}\right)$ satisfying: $c=\frac{M_{N}^{(2)}}{M_{N}^{(1)}}$ $\lim _{N \rightarrow \infty} M_{N}^{(i)}=+\infty, \quad \lim _{N \rightarrow \infty} \frac{M_{N}^{(i)}}{N}=0 \quad$ for $\quad i=1,2, \quad \lim _{N \rightarrow+\infty} \frac{M_{N}^{(2)}}{M_{N}^{(1)}}=0 \quad$ and such that $\left.W\left(M_{N}^{(2)} \theta\right)=W\left(M_{N}^{(1)} \theta\right) \forall \theta \in\right]-\frac{1}{M_{N}^{(1)}}, \frac{1}{M_{N}^{(1)}}[$.

We first show that $f_{N}(\lambda)$ is an asymptotically unbiased estimator of $[\phi(\lambda)]^{\frac{p}{\alpha}}$ for $-\Omega<\lambda<\Omega$ and $\tau \lambda \notin\left\{w_{1}, w_{2}, \cdots, w_{q}\right\}$.

Theorem 3.1
Let $-\Omega<\lambda<\Omega$, such that $\tau \lambda \notin\left\{w_{1}, w_{2}, \cdots, w_{q}\right\}$. Then, $E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}=\mathrm{o}(1)$. If $\phi$ satisfies the hypothesis $|\phi(x)-\phi(y)| \leq \operatorname{cste}|(x-y)|^{-\gamma}$, with $\gamma<2 k \alpha-1$, then,

$$
E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}=\left\{\begin{array}{c}
O\left(\frac{1}{n^{(2 k \alpha-1)}}+\frac{1}{M_{N}^{(1) \gamma}}\right) \quad \text { if } \lambda \neq 0 \\
O\left(\frac{1}{M_{N}^{(1)} n^{(2 k \alpha-1)}}+\frac{1}{M_{N}^{(1) \gamma}}+\frac{1}{n^{2 k \alpha-1}}\right) \quad \text { if } \lambda=0
\end{array}\right.
$$

## Proof

By the definition of the spectral window, we have:

$$
E\left[f_{N}(\lambda)\right]=\int_{R} M_{N}^{(1)} W\left[M_{N}^{(1)}(\lambda-u)\right] E\left[\hat{I}_{N}(u)\right] d u
$$

Let $\quad M_{N}^{(1)}(\lambda-u)=v$ and from (11), we obtain:

$$
\begin{equation*}
E\left[f_{N}(\lambda)\right]=\int_{-1}^{1} W(v)\left[\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)\right]^{\frac{p}{\alpha}} d v \tag{7}
\end{equation*}
$$

Using the fact that $\int_{-1}^{1} W(u) d u=1$ and the inequality (4), we get:

$$
\left|E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}\right| \leq \int_{-1}^{1} W(v)\left|\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v
$$

Since $\frac{p}{\alpha}<1$, we obtain

$$
\begin{aligned}
& \qquad\left|\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} \leq\left|\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}}+\left.\left|\psi_{N, 2}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)\right|\right|^{\frac{p}{\alpha}} \\
& \text { We now examine the limit of } \psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right) \text {. From (3) we get: }
\end{aligned}
$$

$$
\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=\int_{-\infty}^{\infty}\left|H_{N}\left(u-\tau\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)\right)\right|^{\alpha} \phi\left(\frac{u}{\tau}\right) d u .
$$

Let $u-\tau\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=y$, we obtain:

$$
\begin{align*}
\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right) & =\int_{R}\left|H_{N}(y)\right|^{\alpha} \phi\left(\lambda-\frac{v}{M_{N}^{(1)}}+\frac{y}{\tau}\right) d y \\
& =\sum_{j \in Z} \int_{(2 j-1) \pi}^{(2 j+1) \pi}\left|H_{N}(y)\right|^{\alpha} \phi\left(\lambda-\frac{v}{M_{N}^{(1)}}+\frac{y}{\tau}\right) d y . \tag{8}
\end{align*}
$$

Let $y-2 j \pi=s$. Since $\left|H_{N}(.)\right|^{\alpha}$ is $2 \pi-$ periodic function, we get

$$
\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=\sum_{j \in Z} \int_{-\pi}^{+\pi}\left|H_{N}(s)\right|^{\alpha} \phi\left(\lambda-\frac{v}{M_{N}^{(1)}}+\frac{s}{\tau}+\frac{2 \pi}{\tau} j\right) d s
$$

Since the function $\phi$ is uniformly continuous on $[-\Omega, \Omega]$ and the fact that $\left|H_{N}\right|^{\alpha}$ is a kernel, the right hand side of the last equality converges to $\sum_{j \in Z} \phi\left(\lambda+\frac{2 \pi j}{\tau}\right)$. ${ }_{\text {Let }} j$ be an integer such that $-\Omega<\frac{\tau \lambda+2 \pi j}{\tau}<\Omega$. The definition of $\tau$ implies that $|\tau \lambda|<|\tau \Omega|<\pi$. It is easy to
see that $|j|<1$ and then $j=0$. Since $H_{N}^{(1)}$ is a kernel, we obtain that $\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)$
converges to $\phi(\lambda)$. On the other hand,

$$
\psi_{N, 2}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=\sum_{i=1}^{q}\left|H_{N}^{(1)}\left(w_{i}-\tau\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)\right)\right|^{\alpha}
$$

Since $\quad w_{i}$ is different from $\tau \lambda$ and from the lemma 2.2 , we get $\psi_{N, 2}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right) \leq\left[2 \pi\left(\frac{2}{\pi}\right)^{2 k \alpha} n^{2 k \alpha-1}\right]^{-1} \frac{1}{c t e} \sum_{i=1}^{q} c_{i}$, where $\operatorname{cte}=\inf \left|\sin \left(\frac{w_{i}-\tau\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)}{2}\right)\right|$.

Therefore, $\psi_{N, 2}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=O\left(\frac{1}{n^{2 k \alpha-1}}\right)$. Thus, we have $E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}=o(1)$.
The rate of convergence: We assume that the spectral density ${ }^{\phi}$ satisfies the hypothesis $H$. We denote by $F=\left|\operatorname{Bias}\left(f_{N}(\lambda)\right)\right|=\left|E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{p / \alpha}\right|$. It follows that

$$
F \leq \frac{p}{2 \alpha} \int_{-1}^{1} W(v)\left[\left[\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)\right]^{\frac{p}{\alpha}-1}+[\phi(\lambda)]^{\frac{p}{\alpha}-1}\right]\left|\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)\right| d v .
$$

Since $\psi_{N}\left(\lambda-\frac{v}{M_{N}}\right)$ converges to $\phi(\lambda)$, getting the rate of the convergence for $F$ requires to examine the rate of convergence of $\int_{-1}^{1} W(v)\left|\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)\right| d v$ Indeed, from (3), we $\underset{\text { obtain }}{ } \psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)=\int_{-\pi}^{\pi}\left|H_{N}\left(y-\tau \lambda+\frac{\tau v}{M_{N}^{(1)}}\right)\right|^{\alpha} \phi\left(\frac{y}{\tau}\right) d y$.

Denote by $\Delta\left(\psi_{N}, \phi\right)=\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(1)}}\right)-\phi(\lambda)$. Putting $t=-\left(y-\tau \lambda+\frac{\tau v}{M_{N}^{(1)}}\right)$ and using the condition $H$, we get $\left|\Delta\left(\psi_{N}, \phi\right)\right| \leq C_{1} \int_{\tau \lambda-}^{\substack{\tau \lambda-\frac{\tau v}{M_{N}^{(1)}}+\pi}}\left|H_{N}^{(1)}-\pi H_{N}(t)\right|^{\alpha}\left|\frac{v}{M_{N}^{(1)}}+\frac{t}{\tau}\right|^{\gamma} d t$.

It is easy to show that

$$
\begin{aligned}
\int_{-1}^{1} W(v)\left|\Delta\left(\psi_{N}, \phi\right)\right| d v & \leq 2^{\gamma} C_{1}\left|\frac{1}{M_{N}^{(1)}}\right|^{\gamma} \int_{-1}^{1} W(v)|v|^{\gamma} d v \\
& +2^{\gamma} \frac{C_{1}}{\tau^{\gamma}} \int_{-1}^{1} W(v) \int_{\tau_{1} \lambda-\frac{\tau v}{M_{N}^{(1)}-\pi}}^{\tau \lambda-\frac{\tau_{1} v}{M_{N}^{(1)}}+\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t d v
\end{aligned}
$$

The second integral of the right hand side is bounded as follows:

$$
\begin{align*}
\int_{\tau \lambda-\frac{\tau v}{M_{N}^{(1)}}-\pi}^{\tau \lambda-\frac{\tau v}{M_{N}^{(1)}}+\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t & \leq \int_{-|\tau \lambda|-\left|-\frac{\tau v}{M_{N}^{(1)}}\right|-\pi}^{-\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t \\
& +\int_{-\pi}^{\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t \\
& +\int_{\pi}^{|\tau \lambda|+\left.\left|\frac{\tau v}{M_{N}^{(1)} \mid+\pi}\right| H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t} \tag{9}
\end{align*}
$$

The function $\left|H_{N}().\right|$ is even, then the first and the last integrals in the right hand side of (9) are equal. Since $\frac{\tau \nu}{M_{N}}$ converges to zero and $\tau \lambda<\tau \Omega<\pi$, for a large $N$ we have:

$$
\begin{aligned}
\int_{\pi}^{|\tau \lambda|+\left|+\frac{\tau v}{M_{N}^{(1)}}\right|+\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t & \leq(2 \pi)^{\gamma} \int_{\pi}^{|\tau \lambda|+\left|\frac{\tau v}{M_{N}^{(1)}}\right|+\pi}\left|H_{N}(t)\right|^{\alpha} d t \\
& \leq \frac{(2 \pi)^{\gamma}}{B_{\alpha, N}^{\prime}} \frac{|\tau \lambda|+\frac{\tau}{M_{N}^{(1)}}}{\left|\sin \left(\pi+|\tau \lambda|+\frac{\tau}{M_{N}^{(1)}}\right)\right|^{2 k \alpha}}
\end{aligned}
$$

From the lemma 2.1, we obtain,

$$
\int_{\pi}^{\left.|t \lambda|+\frac{1}{M_{N}^{(1)}} \right\rvert\,+\pi}\left|H_{N}(t)\right|^{\alpha}|t|^{\gamma} d t= \begin{cases}\mathrm{O}\left(\frac{1}{n^{2 k \alpha-1}}\right) & \text { if } \lambda \neq 0 \\ \mathrm{O}\left(\frac{1}{M_{N}^{(1)} n^{2 k \alpha-1}}\right) & \text { if } \lambda=0\end{cases}
$$

Thus, the result follows.

Theorem 3.2. Let $\lambda$ a real nuber belonging to $]-\Omega, \Omega\left[\right.$, and $\tau \lambda=w_{i}$. Choose $k$ such that $\lim _{N \rightarrow \infty} \frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}=0$.Then,
i) $\quad E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}=o(1)$
ii) If $\phi$ satisfiies the hypothesis $|\varphi(x)-\varphi(y)| \leq \operatorname{cste}|(x-y)|^{-\gamma},{ }_{\text {with }} \quad \frac{\gamma+1}{2 k}<\alpha<2$, then

$$
\left|E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}\right|=\left\{\begin{array}{lll}
O\left(\frac{1}{\left(M_{N}^{(2)}\right)^{\gamma}}+\frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}\right) & \text { if } & 0<\gamma \leq 1 \\
O\left(\frac{1}{M_{N}^{(2)}}+\frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}\right) & \text { if } & 1<\gamma \leq 2
\end{array}\right.
$$

Proof :
From the defintion of the estimator, we have

$$
\begin{gathered}
E\left[f_{N}(\lambda)\right]=\int_{-\pi}^{\pi} \frac{W_{N}^{(2)}(\lambda-u)-\frac{M_{N}^{(2)}}{M_{N}^{(1)}} W_{N}^{(1)}(\lambda-u)}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}\left[\psi_{N}(u)\right]^{\frac{p}{\alpha}} d u \\
E\left[f_{N}(\lambda)\right]=\int_{-\pi}^{\lambda-\frac{1}{M_{N}^{(1)}}+\int_{\lambda-\frac{1}{M_{N}^{(1)}}}^{\lambda+\frac{1}{M_{N}^{(1)}}}+\int_{\lambda+\frac{1}{M_{N}^{(1)}}}^{\pi}=E_{1}+E_{2}+E_{3}} \\
\lambda-u=v, \quad E_{2}=\int_{-\frac{1}{M_{N}^{(1)}}}^{\frac{1}{M_{N}^{(1)}}} \frac{M_{N}^{(2)}\left[W\left[M_{n}^{(2)} v\right]-W\left[M_{N}^{(1)} v\right]\right]}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}\left[\psi_{N}(\lambda-v)\right]^{\frac{p}{\alpha}} d v . \\
\text { for a large N. }
\end{gathered}
$$

$$
\begin{aligned}
E_{1}= & \frac{M_{N}^{(2)}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{-\pi}^{\lambda-\frac{1}{M_{N}^{(1)}}} W\left[M_{N}^{(2)}(\lambda-u)\right]\left[\psi_{N}(u)\right]^{\frac{p}{\alpha}} d u- \\
& \frac{M_{N}^{(2)}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{-\pi}^{\lambda-\frac{1}{M_{N}^{(1)}}} W\left[M_{N}^{(1)}(\lambda-u)\right]\left[\psi_{N}(u)\right]^{\frac{p}{\alpha}} d u
\end{aligned}
$$

Put $M_{N}^{(2)}(\lambda-u)=v$ in the first integral and put $M_{N}^{(1)}(\lambda-u)=w$, in the second integral and for a large $N$, we have $M_{N}^{(1)}(\lambda+\pi)>1$ et $M_{N}^{(2)}(\lambda+\pi)>1$. As $W$ is null outside of ${ }^{[-1,1]}$, for large $N$, the second integral of $E_{1}$ is zero. Therefore,

$$
\begin{equation*}
E_{1}=\frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}^{1} W(v)\left[\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)\right]^{\frac{p}{\alpha}} d v \tag{10}
\end{equation*}
$$

$E_{3}=\frac{M_{N}^{(2)}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\lambda+\frac{1}{(1)}}^{\pi} W\left[M_{N}^{(2)}(\lambda-u)\right]\left[\psi_{N}(u)\right]^{\frac{p}{\alpha}} d u-\frac{M_{N}^{(2)}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\lambda+\frac{1}{(1)}}^{\pi} W\left[M_{N}^{(1)}(\lambda-u)\right]\left[\psi_{N}(u)\right]^{\frac{p}{\alpha}} d u$

Putting $M_{N}^{(2)}(\lambda-u)=v$ in the first integral and $M_{N}^{(1)}(\lambda-u)=w$ in the second integral, we obtain

$$
E_{3}=\frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{M_{N}^{(2)}(\lambda-\pi)}^{-\frac{M_{N}^{(2)}}{M_{1}^{(1)}}} W(v)\left[\psi_{N}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)\right]^{\frac{p}{\alpha}} d v-\frac{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{M_{N}^{(1)}(\lambda-\pi)}^{-1} W(w)\left[\psi_{N}\left(\lambda-\frac{w}{M_{N}^{(1)}}\right)\right]^{\frac{p}{\alpha}} d w .
$$

For large $N$ we have, $M_{N}^{(2)}(\lambda-\pi)<-1$ and $M_{N}^{(1)}(\lambda-\pi)<-1$.

$$
\begin{equation*}
E_{3}=\frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}^{1} W(v)\left[\psi_{N}\left(\lambda+\frac{v}{M_{N}^{(2)}}\right)\right]^{\frac{p}{\alpha}} d v \tag{11}
\end{equation*}
$$

It is easy to show that for a large $N$

$$
\begin{gather*}
\frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{N}^{1} \frac{M_{N}^{(2)}}{M_{N}^{(1)}} W(v) d v=\frac{1}{2} . \\
\left|E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}\right|=\left|E_{1}+E_{3}-[\phi(\lambda)]^{\frac{p}{\alpha}}\right| \leq\left|E_{1}-\frac{1}{2}[\phi(\lambda)]^{\frac{p}{\alpha}}\right|+\left|E_{3}-\frac{1}{2}[\phi(\lambda)]^{\frac{p}{\alpha}}\right| \tag{12}
\end{gather*}
$$

From (10) and (12), for a large $N$, we have

$$
\left|E_{1}-\frac{1}{2}[\phi(\lambda)]^{\frac{p}{\alpha}}\right| \leq \frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}^{1} W(v)\left|\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)+\psi_{N_{2}}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v
$$

As $\frac{p}{\alpha}<1$, we obtain

$$
\begin{gathered}
\left|E_{1}-\frac{1}{2}[\phi(\lambda)]^{\frac{p}{\alpha}}\right| \leq \frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}} W(v)}^{1}\left|\psi_{N, 1}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v \\
\\
\quad \frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}} \int_{\frac{M_{N}}{(2)}}^{M_{N}^{(1)}} W(v)\left|\psi_{N, 2}\left(\lambda-\frac{v}{M_{N}^{(2)}}\right)\right|^{\frac{p}{\alpha}} d v .}
\end{gathered}
$$

On the other hand,

$$
\int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}^{1} W(v)\left|\psi_{N, 1}\left(\lambda \pm \frac{v}{M_{N}^{(2)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v \leq \int_{0}^{1} W(v)\left|\psi_{N, 1}\left(\lambda \pm \frac{v}{M_{N}^{(2)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v
$$

For all $\lambda$ belonging to $]-\pi, \pi\left[\psi_{N, 1}\left(\lambda \pm \frac{v}{M_{N}^{(2)}}\right) \quad\right.$ converges to $\phi(\lambda)$, uniformly in $v \in[-1,1]$. Therefore, $\frac{1}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{\frac{M_{N}^{(2)}}{M_{N}^{(1)}}}^{1} W(v)\left|\psi_{N, 1}\left(\lambda \pm \frac{v}{M_{N}^{(2)}}\right)-\phi(\lambda)\right|^{\frac{p}{\alpha}} d v$ converge to zero.

$$
\psi_{N, 2}\left(\lambda_{i} \pm \frac{v}{M_{N}^{(2)}}\right) \leq \sum_{m=1}^{q} \frac{a_{m}}{B_{\alpha \neq i}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left(w_{i} \pm \frac{\tau v}{M_{N}^{(2)}}-w_{m}\right)\right]\right|^{2 k \alpha}}
$$

Since $\tau \lambda=w_{i}$,

$$
+\frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left( \pm \frac{\tau v}{M_{N}^{(2)}}\right)\right]\right|^{2 k \alpha}}
$$

For all $m \neq i, \quad \lim _{N \rightarrow+\infty} \sup _{v \in[-1,1]} \frac{1}{\left|\sin \left[\frac{1}{2}\left(w_{i} \pm \frac{\tau v}{M_{N}^{(2)}}-w_{m}\right)\right]\right|^{2 k \alpha}}=\frac{1}{\left\lvert\, \sin \left[\frac{1}{2}\left(w_{i}-w_{m}\right)\right]^{2 k \alpha}\right.}$.
Thus, for large N , we get

$$
\sum_{\substack{m=1 \\ m \geq i}}^{q} \frac{a_{m}}{B_{\alpha, N}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left(w_{i} \pm \frac{\tau v}{M_{N}^{(2)}}-w_{m}\right)\right]\right|^{2 k \alpha}} \leq\left(\varepsilon+\frac{1}{\underset{m \in(1,2, m),(i)}{\inf }\left|\sin \left[\frac{1}{2}\left(w_{i}-w_{m}\right)\right]\right|^{2 k \alpha}}\right) \sum_{\substack{m=1 \\ m \neq i}}^{q} \frac{a_{m}}{B_{\alpha, N}^{\prime}} .
$$

From the lemma 2.1, we obtain

$$
\begin{equation*}
\sum_{m=1}^{q} \frac{a_{m}}{B_{\alpha, i}^{\prime}} \frac{1}{\left.\sin \left[\frac{1}{2}\left(w_{i} \pm \frac{\tau v}{M_{N}^{(2)}}-w_{m}\right)\right]\right|^{2 k \alpha}}=O\left(\frac{1}{n^{2 k \alpha-1}}\right) \tag{13}
\end{equation*}
$$

On the other hand, for large $N$, we have $\sup _{v \in[0,11}\left|\frac{\tau v}{M_{N}^{(2)}}\right|=\frac{\tau}{M_{N}^{(2)}}<\pi$
. Consequently

$$
\frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left( \pm \frac{\tau v}{M_{N}^{(2)}}\right)\right]\right|^{2 k \alpha}} \leq \frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{\pi^{2 k \alpha}}{\left|\frac{\tau v}{M_{N}^{(2)}}\right|^{2 k \alpha}} .
$$

As $\quad \frac{\tau}{M_{N}^{(1)}} \leq \frac{\tau v}{M_{N}^{(2)}} \leq \frac{\tau}{M_{N}^{(2)}}$, we obtain

$$
\frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left( \pm \frac{\tau v}{M_{N}^{(2)}}\right)\right]\right|^{2 k \alpha}} \leq \frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{\pi^{2 k \alpha}}{\left|\frac{\tau}{M_{N}^{(1)}}\right|^{2 k \alpha}}
$$

The lemma 2.2 gives

$$
\begin{equation*}
\frac{a_{i}}{B_{\alpha, N}^{\prime}} \frac{1}{\left|\sin \left[\frac{1}{2}\left(+-\frac{\tau v}{M_{N}^{(2)}}\right)\right]\right|^{2 k \alpha}}=O\left(\frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}\right) \tag{14}
\end{equation*}
$$

Thus, we get

$$
\begin{equation*}
\frac{2^{\frac{p}{\alpha}}}{1-\frac{M_{N}^{(2)}}{M_{N}^{(1)}}} \int_{M_{N}^{(2)}}^{M_{N}^{(1)}} W(v)\left|\psi_{N, 2}\left(\tau \lambda \pm \frac{\tau v}{M_{N}^{(2)}}\right)\right|^{\frac{p}{\alpha}} d v=O\left[\frac{1}{n^{(2 k \alpha-1) \frac{p}{\alpha}}}+\left(\frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}\right)^{\frac{p}{\alpha}}\right] \tag{15}
\end{equation*}
$$

Choosing $M_{N}^{(1)}$ such that $\frac{\left(M_{N}^{(1)}\right)^{2 k \alpha}}{n^{2 k \alpha-1}}$ converges to 0 . For example $M_{N}^{(1)}=n^{b} \quad$ with $0<b<1-\frac{1}{2 k \alpha}$. Thus, $\lim _{N \rightarrow+\infty} E\left[f_{N}(\lambda)\right]-[\phi(\lambda)]^{\frac{p}{\alpha}}=0$.

Theorem 4.2 Let $-\Omega<\lambda<\Omega$ such that $\phi(\lambda)>0$. Then, $\left.\operatorname{var}\left[f_{N}(\lambda)\right)\right]$ converges to zero. If $M_{N}^{(1)}=n_{\text {with }} \frac{1}{2 k^{2} \alpha^{2}}<c<\frac{1}{2}$, then $\left.\quad \operatorname{var}\left[f_{N}(\lambda)\right)\right]=\mathrm{O}\left(\frac{1}{n^{(1-2 c)}}\right)$.

Proof Consider the case where $\tau \lambda \notin\left\{w_{1}, w_{2}, \cdots, w_{q}\right\}$. It is clear that the variance of $f_{N}$ can be written as follows:

$$
\operatorname{var}\left[f_{N}(\lambda)\right]=\int_{R^{2}} W_{N}^{(1)}(\lambda-u) W_{N}^{(1)}\left(\lambda_{1}-u^{\prime}\right) \operatorname{cov}\left[\hat{I}_{N}(u), \hat{I}_{N}\left(u^{\prime}\right)\right] d u d u^{\prime} .
$$

By using the fact that $W$ is zero for $|\lambda|>1$, for large $N$, we get

$$
\operatorname{var}\left[f_{N}(\lambda)\right]=\int_{-1}^{1} \operatorname{cov}\left[\hat{I}_{N}\left(\lambda-\frac{x_{1}}{M_{N}}\right), \hat{I}_{N}\left(\lambda-\frac{x_{1}^{\prime}}{M_{N}}\right)\right] W\left(x_{1}\right) W\left(x_{1}^{\prime}\right) d x_{1} d x_{1}^{\prime} .
$$

We define two subsets of the ${ }^{[-1,1]^{2}}$ by:

$$
\begin{array}{ll}
. & L_{1}=\left\{\left(x_{1}, x_{1}^{\prime}\right) \in[-1,1]^{2} ; \quad\left|x_{1}-x_{1}^{\prime}\right|>\sigma_{N}\right\}, \\
. & L_{2}=\left\{\left(x_{1}, x_{1}^{\prime}\right) \in[-1,1]^{2} ; \quad\left|x_{1}-x_{1}^{\prime}\right| \leq \sigma_{N}\right\},
\end{array}
$$

where $\sigma_{N}$ is a nonnegative real, converging to 0 . We split the integral into an integral over the sub region $L_{2}$ and an integral over $L_{1}: \operatorname{var}\left[f_{N}(\lambda)\right]=\int_{L_{2}}+\int_{L_{1}} \stackrel{\Delta}{=} J_{1}+J_{2}$.

By Cauchy-Schwartz inequality and theorem 3.1, we obtain

$$
J_{1} \leq C \int_{\left|x_{1}-x_{1}\right| \leq \sigma_{N}} W\left(x_{1}\right) W\left(x_{1}^{\prime}\right) d x_{1} d x_{1}^{\prime} \text {. where } C \text { is constant. Thus, we obtain }
$$

$$
\begin{equation*}
J_{1}=O\left(\sigma_{N}\right) \tag{16}
\end{equation*}
$$

It remains to show that $J_{2}$ converges to zero. For simplicity, we define

$$
\lambda_{1}=\lambda-\frac{x_{1}}{M_{N}^{(1)}} ; \quad \lambda_{2}=\lambda-\frac{x_{1}^{\prime}}{M_{N}^{(1)}}, \quad C(\lambda)=\operatorname{cov}\left[\hat{I}_{N}\left(\lambda-\frac{x_{1}}{M_{N}^{(1)}}\right), \hat{I}_{N}\left(\lambda-\frac{x_{1}^{\prime}}{M_{N}^{(1)}}\right)\right] .
$$

We first show that $C(\lambda)$ converges to zero uniformly in $x_{1}, x_{1}^{\prime} \in[-1,1]$. Indeed, from lemma 2.3, we have

$$
E \hat{I}_{N}(v)-\hat{I}_{N}(v)=F_{p, \alpha}^{-1}\left[C_{\alpha}\right]^{-p / \alpha} \int_{-\infty}^{\infty} \frac{\operatorname{Re}\left(e^{i u I_{N}(v)}\right)-e^{-C_{\alpha}|u|^{\alpha} \psi_{N}(v)}}{|u|^{1+p}} d u
$$

Thus, the expression of the covariance becomes

$$
\begin{aligned}
C(\lambda)= & F_{p, \alpha}^{-2} C_{\alpha}^{-\frac{2 p}{\alpha}} \int_{R^{2}} E\left[\prod_{k=1}^{2} \cos \left(u_{k} I_{N}\left(\lambda_{k}\right)\right)\right] \\
& -\exp \left\{-C_{\alpha} \sum_{k=1}^{2}\left|u_{k}\right|^{\alpha} \psi_{N}\left(\lambda_{k}\right)\right\} \frac{d u_{1} d u_{2}}{\left|u_{1} u_{2}\right|^{1+p}}
\end{aligned}
$$

The following equality : $2 \cos x \cos y=\cos (x+y)+\cos (x-y)$, implies that

$$
\begin{aligned}
E\left[\prod_{k=1}^{2}\left(\cos u_{k} I_{N}\left(\lambda_{k}\right)\right)\right] & =\frac{1}{2} \exp \left[-C_{\alpha} \int\left|(\tau)^{\frac{1}{\alpha}} \sum_{k=1}^{2} u_{k} H_{N}\left(\tau \lambda_{k}-\tau v\right)\right|^{\alpha} d \mu(v)\right] \\
& +\frac{1}{2} \exp \left[-C_{\alpha} \int\left|(\tau)^{\frac{1}{\alpha}} \sum_{k=1}^{2}(-1)^{k+1} u_{k} H_{N}\left(\tau \lambda_{k}-\tau v\right)\right|^{\alpha} d \mu(v)\right]
\end{aligned}
$$

By substituting the expression for $C(\lambda)$ and changing the variable ${ }^{u_{2}}$ to $\left(-u_{2}\right)$ in the second term, we obtain

$$
\begin{equation*}
C(\lambda)=F_{p, \alpha}^{-2} C_{\alpha}^{-\frac{2 p}{\alpha}} \int_{R^{2}}\left(e^{-K}-e^{-K^{\prime}}\right) \frac{d u_{1} d u_{2}}{\left|u_{1} u_{2}\right|^{1+p}}, \tag{17}
\end{equation*}
$$

$\underset{\text { where }}{ } K=C_{\alpha} \int_{R}\left|(\tau)^{\frac{1}{\alpha}} \sum_{k=1}^{2} u_{k} H_{N}\left(\tau \lambda_{k}-\tau v\right)\right|^{\alpha} d \xi(v)$ and

$$
K^{\prime}=C_{\alpha} \tau \sum_{k=1}^{2}\left|u_{k}\right|^{\alpha} \int_{R}\left|H_{N}\left(\tau \lambda_{k}-\tau v\right)\right|^{\alpha} \xi(v) d v
$$

Since $\quad K, K^{\prime}>0, \quad\left|e^{-K}-e^{-K^{\prime}}\right| \leq\left|K-K^{\prime}\right| \exp \left\{\left|K-K^{\prime}\right|-K^{\prime}\right\} \quad$, we obtain: $\left|K-K^{\prime}\right| \leq 2 C_{\alpha} \tau|u|^{\frac{\alpha}{2}} Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)$,
where
$Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)=\int_{-\Omega}^{\Omega}\left|H_{N}\left(\tau \lambda_{1}-\tau u\right)\right|^{\frac{\alpha}{2}}\left|H_{N}\left(\tau \lambda_{2}-\tau u\right)\right|^{\frac{\alpha}{2}} d \xi(u)$
Now, let us show that $Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)$ converges to zero. Indeed, since ${ }^{\phi}$ is bounded on $[-\Omega, \Omega]$, we have

$$
\begin{align*}
Q_{N}\left(\lambda_{1} ; \lambda_{2}\right) & \leq \sup (\phi) \int_{-\Omega}^{\Omega}\left|H_{N}\left(\tau \lambda_{1}-\tau u\right) H_{N}\left(\tau \lambda_{2}-\tau u\right)\right|^{\frac{\alpha}{2}} d u \\
& +\sum_{i=1}^{q} c_{i}\left|H_{N}\left(\tau \lambda_{1}-\tau w_{i}\right) H_{N}\left(\tau \lambda_{2}-\tau w_{i}\right)\right|^{\frac{\alpha}{2}} \tag{18}
\end{align*}
$$

From the definition of $H_{N}$, we can write

$$
\int_{-\Omega}^{\Omega}\left|H_{N}\left(\tau \lambda_{1}-\tau v\right) H_{N}\left(\tau \lambda_{2}-\tau v\right)\right|^{\frac{\alpha}{2}} d v=\left.\int_{-\Omega}^{\Omega} \frac{1}{B_{\alpha, N}^{\prime}}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|\right|^{k \alpha} d v
$$

## a) First step:

We show that the denominators of the first and second terms under the last integral do not vanish for the same $v$, so we suppose that $v_{\text {exists, belonging to }}[-\Omega, \Omega]$ and $z, z^{\prime} \in Z$ such as: $\tau \lambda_{1}-\tau v=2 z \pi$ and $\tau \lambda_{2}-\tau v=2 z^{\prime} \pi$. Since $\lambda_{1} \neq \lambda_{2}$, then $z$ and $z^{\prime}$ are different. Therefore, $\quad z-z^{\prime}=\frac{\tau}{2 \pi}\left(\lambda_{1}-\lambda_{2}\right)$. Hence, $\quad\left|z-z^{\prime}\right|=\frac{1}{w}\left|\lambda_{1}-\lambda_{2}\right|$. As $\quad \lim _{N \rightarrow \infty}\left|\lambda_{1}-\lambda_{2}\right|=0$, consequently, for a large $N_{\text {we get: }}|z-z|<\frac{1}{2}$.

Thus, we obtain a contradiction with the fact that $z$ and $z$ are different integers.
b) Second step

We assume there exist $q$ points, $\quad V_{1}, V_{2}, \cdots, V_{q} \in[-\Omega, \Omega]$ such as:
for $j=1,2, \cdots, q \quad \tau \lambda_{1}-\tau V_{j} \in 2 \pi Z \quad$ then $\frac{\lambda_{1}}{w}-\frac{V_{j}}{w} \in Z \quad$, and we assume that there exist $q^{\prime}$ points $V_{1}^{\prime}, V_{2}^{\prime}, \cdots, V_{q^{\prime}}^{\prime} \in[-\Omega, \Omega]$ such that for $i=1,2, \cdots, q^{\prime} \quad \frac{\lambda_{2}}{w}-\frac{V_{i}^{\prime}}{w} \in Z$. Showing that, $\left|V_{j}\right| \neq \Omega, \quad\left|V_{i}^{\prime}\right| \neq \Omega \quad$ for $\quad 1 \leq j \leq q \quad$ and $\quad 1 \leq i \leq q^{\prime} . \quad$ Indeed, $\quad-1<\frac{\lambda-\Omega}{w}<0 \quad$ and $0<\frac{\lambda+\Omega}{w}<1$ because $w>2 \Omega$. Hence $\frac{\lambda-\Omega}{w} \notin Z \quad$ and $\frac{\lambda+\Omega}{w} \notin Z$. For a large $N$, we get that $\left[\frac{\lambda-\Omega}{w}\right]_{E}<\frac{\lambda_{1}}{w}-\frac{\Omega}{w}<1+\left[\frac{\lambda-\Omega}{w}\right]_{E}, \quad$ where $[x]_{E}$ is the integer part of $x$. Hence, $\frac{\lambda_{1}}{w}-\frac{\Omega}{w} \notin Z \quad$. In the same manner, we show that $\frac{\lambda_{1}}{w}+\frac{\Omega}{w} \notin Z$ . Similarly, it can be shown that $\frac{\lambda_{2}+\Omega}{w} \notin Z$. Thus, $\left|V_{j}\right| \neq \Omega$ and $\left|V_{i}^{\prime}\right| \neq \Omega$.

## c) Third step

We classify $V_{j}$ and $V_{i}^{\prime}$ by increasing order: $\quad-\Omega<V_{j_{1}}<V_{j_{2}}<\ldots<V_{j^{\prime}, q^{\prime}}<\Omega$, and we write the integral in the following manner:

$$
\int_{-\Omega}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau_{1} v\right)\right]}\right|^{k \alpha} d v=I_{1}+\sum_{i=1}^{q+q^{\prime}} I_{2, i}+\sum_{i=1}^{q+q^{\prime}-1} I_{3, i}+I_{4}
$$

where

$$
\begin{aligned}
& I_{1}=\int_{-\Omega_{1}}^{V_{j_{1}}-\delta(N)}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \\
& I_{2, i}=\int_{V_{j_{i}}-\delta(N)}^{V_{j_{j}}+\delta(N)}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \\
& I_{3, i}=\int_{V_{j_{i}}+\delta(N)}^{V_{j_{i+1}}-\delta(N)}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \\
& I_{4}=\int_{V_{j}{ }_{q+q}+\delta(N)}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau_{1} \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v
\end{aligned}
$$

where $\delta(N)$ is a nonnegative real number converging to zero and satisfying:
$-\Omega<V_{j_{1}}-\delta(N)<V_{j_{1}}+\delta(N)<V_{j_{2}}-\delta(N)<V_{j_{2}}+\delta(N)<\ldots \ldots<V_{j_{q+q^{\prime}}}-\delta(N)<V_{j^{\prime}+q^{\prime}}+\delta(N)<\Omega$,
$\underset{\text { and }}{ } \delta(N)<\left|\frac{\lambda_{1}-\lambda_{2}}{2}\right|$.
First, we show that the first integral converges to zero. We know that for a large $N$, we have $\lambda_{1}<\Omega$. For simplicity without loss of generality, we assume that for all ${ }^{i}, \tau \lambda_{1}-\tau V_{j_{i}}=2 k_{i} \pi$ with $k_{i} \in Z$. Similarly, we can show the result if we rather assume that $\tau \lambda_{2}-\tau V_{j_{i}}=2 k_{i} \pi$ with $k_{i} \in Z$.Since there is no $v$ between $-\Omega$ and $V_{j_{1}}-\delta(N)$ on which the denominators are vanishing,

$$
\begin{aligned}
& I_{1} \leq \frac{V_{j_{1}}-\delta(N)+\Omega}{\inf \left[\left|\sin \frac{\tau \delta(N)}{2}\right|^{k \alpha},\left|\sin \frac{\tau\left(\lambda_{1}+\Omega\right)}{2}\right|^{k \alpha}\right] \inf \left[\left|\sin \frac{\tau\left(\lambda_{2}-V_{j_{1}}+\delta(N)\right)}{2}\right|^{k \alpha},\left|\sin \frac{\tau\left(\lambda_{2}+\Omega\right)}{2}\right|^{k \alpha}\right]} \\
& \left|\sin \frac{\tau\left(\lambda_{2}-V_{j_{1}}+\delta(N)\right)}{2}\right|^{k \alpha}=\left|\sin \frac{\tau\left|\lambda_{2}-\lambda_{1}+\delta(N)\right|}{2}\right|^{k \alpha} \quad \text {. For a large } N \text {, we have }
\end{aligned}
$$

$\frac{\tau\left|\lambda_{2}-\lambda_{1}+\delta(N)\right|}{2} \leq \tau \Omega+\frac{\tau \delta(N)}{2}<\pi-\frac{\tau \delta(N)}{2}$. On the other hand, two cases are possible:

1. $\lambda_{2}-\lambda_{1}>0$, then we have $\left|\lambda_{2}-\lambda_{1}+\delta(N)\right|=\lambda_{2}-\lambda_{1}+\delta(N)>\delta(N)$
2. $\lambda_{2}-\lambda_{1}<0$, since $\left|\lambda_{2}-\lambda_{1}\right|>2 \delta(N)$, we have
$\left|\lambda_{2}-\lambda_{1}+\delta(N)\right|=\lambda_{1}-\lambda_{2}-\delta(N)>\delta(N)$.
Therefore, $\quad \frac{\tau \delta(N)}{2}<\frac{\tau\left|\lambda_{2}-\lambda_{1}+\delta(N)\right|}{2}<\pi-\frac{\tau \delta(N)}{2}$. For a large $N$, we have $\tau \delta(N)<2 \pi-\tau\left(\lambda_{1}+\Omega\right)$ and $\tau \delta(N)<2 \pi-\tau\left(\lambda_{2}+\Omega\right)$. Then,

$$
\begin{aligned}
& \frac{\tau \delta(N)}{2}<\frac{\tau\left(\lambda_{1}+\Omega\right)}{2}<\pi-\frac{\tau \delta(N)}{2} \text { and } \frac{\tau \delta(N)}{2}<\frac{\tau\left(\lambda_{2}+\Omega\right)}{2}<\pi-\frac{\tau \delta(N)}{2} . \text { Consequently, } \\
& I_{1} \leq \frac{V_{j_{1}}-\delta(N)+\Omega}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{2 k \alpha}}
\end{aligned}
$$

For the integral $I_{2, i}$, we bound the first fraction under integral by $n^{k \alpha}$ :

$$
I_{2, i} \leq n^{k \alpha} \int_{V_{j_{i}}-\delta(N)}^{V_{j^{\prime}}+\delta(N)} \frac{1}{\left\lvert\, \sin \left[\left.\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right|^{\mid k \alpha}\right.\right.} d v .
$$

. By substituting for $V_{j_{i}}$ in the last inequality and putting $v=u-\frac{2 k \pi}{\tau}$, we get

$$
I_{2, i} \leq n^{k \alpha} \int_{\lambda_{1}-\delta(N)}^{\lambda_{1}+\delta(N)} \frac{1}{\left\lvert\, \sin \left[\left.\frac{1}{2}\left(\tau \lambda_{2}-\tau u\right)\right|^{k \alpha}\right.\right.} d u
$$

$$
\text { Since }\left|\lambda_{1}-u\right|<\delta(N) \text {, it is easy to note that: }
$$

$$
\left|\lambda_{2}-u\right| \geqslant\left|\lambda_{2}-\lambda_{1}\right|-\left|\lambda_{1}-u\right| \geqslant\left|\lambda_{2}-\lambda_{1}\right|-\delta(N)>\frac{\left|\lambda_{2}-\lambda_{1}\right|}{2}
$$

Since $\delta(N)$ converges to zero, for a large $N$ we have $\delta(N)<\frac{2}{\tau}\left(\pi-\frac{\tau}{2}\left|\lambda_{2}-\lambda_{1}\right|\right)$, therefore $0<\tau \frac{\left|\lambda_{2}-\lambda_{1}\right|}{4}<\tau \frac{\left|\lambda_{2}-u\right|}{2}<\tau \frac{\left|\lambda_{2}-\lambda_{1}\right|+\delta(N)}{2}<\pi$.

$$
I_{2, i} \leq \frac{2 \delta(N) n^{k \alpha}}{\inf \left[\sin \frac{\left|\lambda_{1}-\lambda_{2}\right| \tau}{4} ; \sin \frac{\left|\lambda_{1}-\lambda_{2}\right| \tau+\delta(N) \tau}{2}\right]}
$$

where $\left|\lambda_{2}-\lambda_{1}\right|=\left|\frac{x_{2}-x_{1}}{M_{N}}\right|$. Then, for a large $N$, we have $\frac{3 \tau}{2}\left|\lambda_{2}-\lambda_{1}\right|<2 \pi$.
Therefore, $\tau \frac{\left|\lambda_{2}-\lambda_{1}\right|}{4}<\frac{\tau\left|\lambda_{2}-\lambda_{1}\right|+\tau \delta(N)}{2}<\pi-\tau \frac{\left|\lambda_{2}-\lambda_{1}\right|}{4}$.

Thus, we bound the integral as follows:

$$
I_{2, i} \leq \frac{2 \delta(N) n^{k \alpha}}{\left|\sin \frac{\tau\left|\lambda_{2}-\lambda_{1}\right| \mid}{4}\right|^{k \alpha}}
$$

Since there is no $v$ between $V_{j_{i}}+\delta(N)$ and $V_{j_{i+1}}-\delta(N)$ on which the denominators are vanishing, we get:

$$
I_{3, i} \leq \frac{V_{j_{i+1}}-V_{j_{i}}-2 \delta(N)}{A \times B},
$$

where

$$
\begin{aligned}
& A=\inf \left(\left|\sin \frac{\tau \lambda_{1}-\tau V_{j_{i}}-\tau \delta(N)}{2}\right|^{k \alpha},\left|\sin \frac{\tau \lambda_{1}-\tau V_{j_{i+1}}+\tau \delta(N)}{2}\right|^{k \alpha}\right) \\
& B=\inf \left(\left|\sin \frac{\tau \lambda_{2}-\tau V_{j_{i}}-\tau \delta(N)}{2}\right|^{k \alpha},\left|\sin \frac{\tau \lambda_{2}-\tau V_{j_{i+1}}+\tau \delta(N)}{2}\right|^{k \alpha}\right)
\end{aligned}
$$

It follows from the hypothesis on $\delta(N)$ that
$\frac{\tau \delta(N)}{2}<\frac{\tau\left|\lambda_{1}-\lambda_{2}\right|}{2}-\frac{\tau|\delta(N)|}{2}<\frac{\tau\left|\lambda_{1}-\lambda_{2}-\delta(N)\right|}{2}<\frac{\tau\left|\lambda_{1}-\lambda_{2}\right|}{2}+\frac{\tau|\delta(N)|}{2}<\pi-\frac{\tau \delta(N)}{2}$.

Thus, by using the definition of ${ }^{V_{j_{i}}}$, we obtain
$\frac{1}{\left|\sin \frac{\tau \lambda_{2}-\tau V_{j_{i}}-\tau \delta(N)}{2}\right|^{k \alpha}}<\frac{1}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{k \alpha}}$.
We use the same way for bounding the other terms

$$
I_{3, i} \leq \frac{V_{j_{i+1}}-V_{j_{i}}-2 \delta(N)}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{2 k \alpha}}
$$ Similarly, since there is no $v$ between $V_{j_{q+q^{\prime}}}+\delta(N)$

and we get : hing, we can show that:
$I_{4} \leq \frac{V_{j_{i+1}}-V_{j_{i}}-2 \delta(N)}{E \times F}, \quad$ where

$$
\begin{aligned}
& E=\inf \left(\left|\sin \frac{\tau \lambda_{1}-\tau \Omega}{2}\right|^{k \alpha},\left|\sin \frac{\tau \lambda_{1}-\tau V_{j_{q+q^{\prime}}}-\tau \delta(N)}{2}\right|^{k \alpha}\right) \\
& F=\inf \left(\left|\sin \frac{\tau \lambda_{2}-\tau \Omega}{2}\right|^{k \alpha},\left|\sin \frac{\tau \lambda_{2}-\tau V_{j_{q+q^{\prime}}}-\tau \delta(N)}{2}\right|^{k \alpha}\right)
\end{aligned}
$$

Since $\delta(N)$ converges to zero, for a large $N$, we have $\delta(N)<\frac{2}{\tau}\left(\pi-\frac{\tau}{2}\left|\lambda_{1}-\Omega\right|\right)$, and $\delta(N)<\frac{2}{\tau}\left(\pi-\frac{\tau}{2}\left|\lambda_{1}-\lambda_{2}\right|\right)$. It follows that: $\quad I_{4} \leq \frac{\Omega-V_{j_{q+q^{\prime}}}-\delta(N)}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{2 k \alpha}}$.

We recapitulate, from the previous increases, we obtain

$$
\begin{aligned}
\left.\int_{-\Omega}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha} \right\rvert\, & \left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha}
\end{aligned} \quad d v \leq \frac{\Omega+V_{j_{1}}-\delta(N)}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{2 k \alpha}}+\sum_{i=1}^{q+q^{\prime}} \frac{2 n^{k \alpha} \delta(N)}{\left|\sin \frac{\tau\left|\lambda_{2}-\lambda_{1}\right|}{4}\right|^{k \alpha}}
$$

After simplification, we have

$$
\int_{-\Omega}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \leq \frac{2 \Omega-2\left(q+q^{\prime}+1\right) \delta(N)}{\left|\sin \frac{\tau \delta(N)}{2}\right|^{2 k \alpha}}+\frac{2 n^{k \alpha} \delta(N)\left(q+q^{\prime}\right)}{\left|\sin \frac{\left.\tau\left|\lambda_{2}-\lambda_{1}\right|\right|^{k \alpha}}{4}\right|^{k} .}
$$

By bounding the first term on the right hand side of the last inequality and using the following inequality $\left|\sin \frac{x}{2}\right| \geq \frac{|x|}{\pi}$, we get

$$
\int_{-\Omega}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \leq \frac{2 \Omega \pi^{2 k \alpha}}{(\tau \delta(N))^{2 k \alpha}}+\frac{2 n^{k \alpha} \delta(N)\left(q+q^{\prime}\right)(2 \pi)^{k \alpha}}{\left(\frac{\tau\left|x_{2}-x_{1}\right|}{M_{N}}\right)^{k \alpha}}
$$

It follows from the lemma 2.1

$$
\begin{align*}
\frac{1}{B_{\alpha, N}^{\prime}} \int_{-\Omega}^{\Omega}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{1}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{1}-\tau v\right)\right.}\right|^{k \alpha}\left|\frac{\sin \left[\frac{n}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}{\sin \left[\frac{1}{2}\left(\tau \lambda_{2}-\tau v\right)\right]}\right|^{k \alpha} d v \leq \\
\quad \frac{1}{\pi}\left(\frac{\pi}{2}\right)^{2 k \alpha}\left(\frac{2 \Omega \pi^{2 k \alpha}}{n^{2 k \alpha-1}(\tau \delta(N))^{2 k \alpha}}+\frac{2 \delta(N)\left(q+q^{\prime}\right)(2 \pi)^{k \alpha}}{n^{k \alpha-1}\left(\frac{\tau \sigma_{N}}{M_{N}}\right)^{k \alpha}}\right) . \tag{19}
\end{align*}
$$

In order to obtain the convergence of the last expression to zero, we choose $\delta(N)=n^{-\beta}, \quad \beta>0$, such as

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{n^{2 k \alpha \beta}}{n^{2 k \alpha-1}}=0 \quad \text { and } \quad \lim _{n \rightarrow \infty} \frac{1}{n^{\beta+k \alpha-1}\left(\frac{\sigma_{N}}{M_{N}}\right)^{k \alpha}}=0 . \tag{20}
\end{equation*}
$$

Thus, from (18) $Q_{N}$ converges to zero. On the other hand,

$$
C(\lambda) \leq F_{p, \alpha}^{-2} C_{\alpha}^{-\frac{2 p}{\alpha}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty}\left|K-K^{\prime}\right| e^{\left|K-K^{\prime}\right|-K^{\prime}} \frac{d u_{1} d u_{2}}{\left|u_{1} u_{2}\right|^{1+p}},
$$

where $\left|K-K^{\prime}\right|-K^{\prime} \leq-C_{\alpha} \sum_{k^{\prime}=1}^{2}\left|u_{k^{\prime}}\right|^{\alpha}\left|\psi_{N}\left(\lambda_{k^{\prime}}\right)-\tau Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)\right|$.

We denote by : $\delta_{\left(N, k^{\prime}\right)}=\psi_{N}\left(\lambda_{k^{\prime}}\right)-\tau Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)$. It follows from (18) and (20) that $\delta_{\left(N, k^{\prime}\right)}$ converges to $\phi(\lambda)$. Hence,

$$
C(\lambda) \leq F_{p, a}^{-2} C_{\alpha}^{-\frac{2 p}{\alpha}} 2 \tau C_{\alpha} Q_{N}\left(\lambda_{1} ; \lambda_{2}\right) 4 \prod_{k^{\prime}=1}^{2} \int_{0}^{\infty} \exp \left[-C_{\alpha}\left(u_{k^{\prime}}\right)^{\alpha}\left|\delta_{\left(N, k^{\prime}\right)}\right|\right] \frac{d u_{k^{\prime}}}{\left(u_{k^{\prime}}\right)^{1+p-\frac{\alpha}{2}}} .
$$

Putting $\quad u_{k^{\prime}}\left(\delta_{N, k^{\prime}}\right)^{\frac{1}{\alpha}}=v$, we obtain

$$
\begin{equation*}
C(\lambda) \leq F_{p, \alpha}^{-2} C_{\alpha}^{-\frac{2 p}{\alpha}} 2 \tau C_{\alpha} \frac{Q_{N}\left(\lambda_{1} ; \lambda_{2}\right)}{\left[\delta_{(N, 1)} \delta_{(N, 2)}\right]^{\frac{1}{2}-\frac{p}{\alpha}}}\left(\int_{-\infty}^{\infty} \frac{e^{-C_{\alpha}|v|^{\alpha}}}{|v|^{1+p-\frac{\alpha}{2}}} d v\right)^{2} \tag{20}
\end{equation*}
$$

Since $\phi(\lambda)>0, C(\lambda)$ converges uniformly in $x_{1}, x_{1}^{\prime} \in[-1,1]$ to zero. From (20), we obtain
$J_{2}=O\left(\frac{1}{n^{2 k \alpha(1-\beta)-1}}+\frac{1}{n^{\beta+k \alpha-1}\left(\frac{\sigma_{N}}{M_{N}}\right)^{k \alpha}}\right)$. is an asymptotically unbiased and consistent estimator.

## 4. CONCLUSIONS

In this paper, we estimate the spectral density of mixed stable process with continuous time when the process is observed at discrete instants. The aliasing phenomenon is avoided by assuming that the spectral is a compact support. This work can be applied in various fields. For example:

- The study of soil cracking where the observed signal is the resistance of the soil. This resistance can have random jumps that are due to the presence of some stones. Thus, the spectral measurement will be composed of two parts, one continuous and the other discrete corresponding to the resistance jumps encountered during the measurement.
- The growth of fruit on a tree can be considered as a continuous distribution, and when there is a fall of a fruit, the other fruits remaining on the tree absorb more energy and their growth will have a jump in value.

The perspective of this work is to optimize the smoothing parameters to have a better rate of convergence. For this purpose, the cross-validation method will be the most appropriate tool.

This work can also be completed by considering a more general case when we observe the process with random errors. In this case, we will use the deconvolution methods, which have proved their efficiency in the presence of random errors.

It would be interesting to give an estimator of the mode of the spectral density representing the frequency where the spectral density reaches the maximum of energy. For that, we must estimate the derivative of the spectral density function.

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#### Abstract

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